

Neutron-Antineutron Oscillation Experiments: What Have We Learned at the Workshop?

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Indiana University/CEEM
Project X Workshop

Why $\Delta B=2$? (theory/phenomenology)

Neutron-antineutron oscillations in nuclei: theory and
experiment

Free neutron oscillations: experimental requirements
@Project X

Thanks to co-conveners: Chris Quigg (FNAL), Albert Young (NC State)

Neutron-Antineutron Oscillations: Speaker List (from Germany, Georgia, India, Japan, US)

Speaker	Subject
R. Mohapathra, Maryland	theory/phenomonology
M. Snow, Indiana	various
G. Greene, ORNL/Tennessee	R&D needs
I. Gogoladze, Bartol/Delaware	theory/phenomonology
M. Chen, Irvine	leptogenesis
K. Babu, Oklahoma State	theory/phenomonology
M. Stavenga, FNAL	theory
M. Buchoff, LLNL	theory/lattice
E. Kearns, Boston	experiment/ $\bar{n}n$ in nuclei
A. Vainshtein, Minnesota	theory/ $\bar{n}n$ in nuclei
Y. Kamyshev, Tennessee	experiment options
R. Tayloe, Indiana	detectors
K. Ganezer, CSUDH	$\bar{n}n$ in nuclei
D. Dubbers, Heidelberg	ILL experiment
T. Gabriel, ORNL/Tennessee	SNS 1MW target
G. Muhrer, LANL	1MW target/moderator design
H. Shimizu, Nagoya	neutron supermirror optics
C-Y Liu, Indiana (also for D. Baxter, Indiana)	moderator experiments/simulations
S. Banerjee, Tata Institute	detectors

Neutron-Antineutron Oscillations: Formalism

$$\Psi = \begin{pmatrix} n \\ \bar{n} \end{pmatrix} \quad \text{n-nbar state vector}$$

$\alpha \neq 0$ allows oscillations

$$H = \begin{pmatrix} E_n & \alpha \\ \alpha & E_{\bar{n}} \end{pmatrix} \quad \text{Hamiltonian of n-nbar system}$$

$$E_n = m_n + \frac{p^2}{2m_n} + U_n \quad ; \quad E_{\bar{n}} = m_{\bar{n}} + \frac{p^2}{2m_{\bar{n}}} + U_{\bar{n}}$$

Note :

- α real (assuming T)
- $m_n = m_{\bar{n}}$ (assuming CPT)
- $U_n \neq U_{\bar{n}}$ in matter and in external B [$\mu(\bar{n}) = -\mu(n)$ from CPT]

Neutron-Antineutron transition probability

$$\text{For } H = \begin{pmatrix} E + V & \alpha \\ \alpha & E - V \end{pmatrix} \quad P_{n \rightarrow \bar{n}}(t) = \frac{\alpha^2}{\alpha^2 + V^2} \times \sin^2 \left[\frac{\sqrt{\alpha^2 + V^2}}{\hbar} t \right]$$

where V is the potential difference for neutron and anti-neutron.

Present limit on $\alpha \leq 10^{-23} \text{ eV}$

Contributions to V :

$\langle V_{\text{matter}} \rangle \sim 100 \text{ neV}$, proportional to density

$\langle V_{\text{mag}} \rangle = \mu B$, $\sim 60 \text{ neV/Tesla}$; $B \sim 10 \text{ nT} \rightarrow V_{\text{mag}} \sim 10^{-15} \text{ eV}$

$\langle V_{\text{matter}} \rangle$, $\langle V_{\text{mag}} \rangle$ both $\gg \alpha$

$$\text{For } \left[\frac{\sqrt{\alpha^2 + V^2}}{\hbar} t \right] \ll 1 \text{ ("quasifree condition")} \quad P_{n \rightarrow \bar{n}} = \left(\frac{\alpha}{\hbar} \times t \right)^2 = \left(\frac{t}{\tau_{n\bar{n}}} \right)^2$$

Figure of merit = NT^2 $N = \# \text{neutrons}$, $T = \text{"quasifree" observation time}$

How to Search for N-Nbar Oscillations

Figure of merit for probability:

$$NT^2$$

N=total # of free neutrons observed

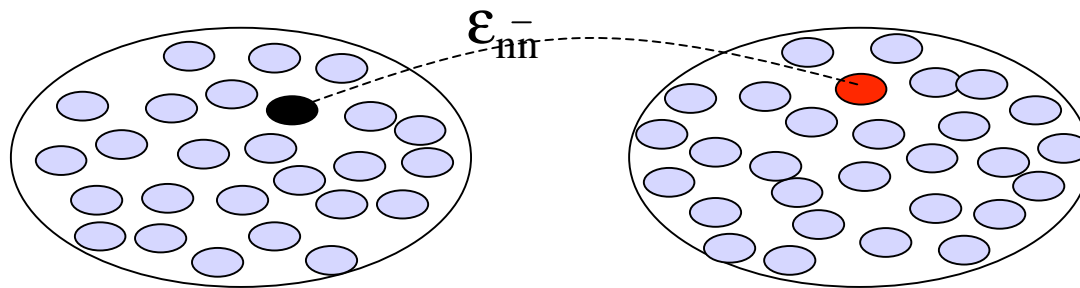
T= observation time per neutron while in “quasifree” condition

When neutrons are in matter or in nucleus, n-nbar potential difference is large->quasifree observation time is short

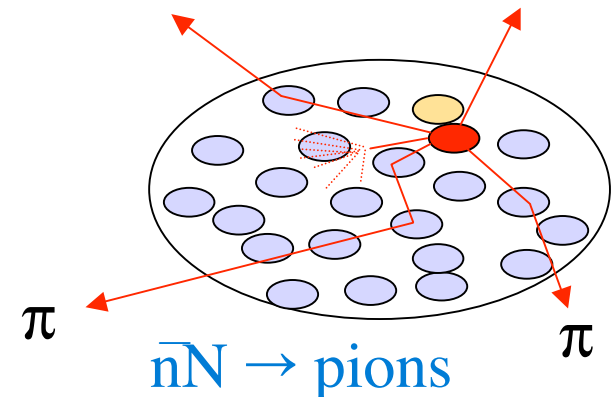
B field must be suppressed to maintain quasifree condition due to opposite magnetic moments for neutron and antineutron

(1) n-nbar transitions in nuclei in underground detectors

(2) Cold and Ultracold neutrons



Nucleus $A \rightarrow A^* + \bar{n}$





Why is it important to search for $NN\bar{b}$?

- Many reasons to believe that **baryon number (B) is not a good symmetry of nature** :
Sphalerons in SM , GUTs, origin of matter etc.
- If B is violated, important to determine the selection rules: $B=1$ (p-decay) or $B=2$ ($NN\bar{b}$) ?
 - i) What is the scale at which B- symmetry is broken ?
 $NN\bar{b} \rightarrow$ lower scale physics than usual p-decay
 - ii) $NN\bar{b}$ oscillation intimately connected to neutrino mass physics when combined with quark-lepton unification



Questions for N-N-bar oscillation

- Are there decent (**predictive?**) theories explaining small neutrino masses which give observable N-N-bar oscillation ?
- Implications of observable N-N-bar for cosmology i.e. does it affect conventional explanations of origin of matter/can it explain itself ?
- Two examples of models for NNbar:

- (i) **TeV scale Seesaw + Quark-Lepton unif.**
- (ii) **SO(10) GUT scale seesaw + TeV sextets**

New Particles at LHC:

Color sextet scalars Δ_{qq}

- TeVColor sextets are an inherent part of both models ;
Can be searched at LHC:

(I) **Single production:** $ud \rightarrow \Delta_{ud} \rightarrow tj$

xsection calculated in (RNM, Okada, Yu' 07;) resonance peaks above SM background- decay to tj;

- **Important LHC signature:** $\sigma(tt) > \sigma(\bar{t}\bar{t})$

(II) **Drell-Yan pair production** $q\bar{q} \rightarrow G \rightarrow \Delta_{ud}\bar{\Delta}_{ud}$

- Leads to $tj\bar{t}j$ final states: **LHC reach < TeV**

Origin of matter and neutron oscillation

■ Current scenarios:

- (i) Leptogenesis; Related to seesaw; but hard to test !
- (ii) Electroweak baryogenesis :

$M_{\text{higgs}} < 127 \text{ GeV}; \quad m_{\tilde{t}} \leq 120 \text{ GeV}$ (puts MSSM under tension)

■ New scenarios: (Babu's talk)

- (iii) Post sphaleron Baryogenesis
 - (iv) GUT baryogenesis
- } both connected to $NN\bar{b}$ osc.

- Non-observation of $NN\bar{b}$ upto 10^{11} sec. will rule out simple models for PSB as well as the particular $SO(10)$ model.

Summary and Conclusions

Post-sphaleron baryogenesis predicts observable $n - \bar{n}$ oscillations

Colored scalars at TeV scale should be accessible to LHC

New GUT scale $(B - L)$ -genesis proposed which is sphaleron-proof

Both models predict

$$\tau(n - \bar{n}) \sim (10^9 - 10^{11}) \text{ sec}$$

$n - \bar{n}$ oscillation experiments can probe a class of theories which explains the origin of matter in the universe

Conclusions

- origin of matter: one of the great mysteries in particle physics and cosmology
- leptogenesis: an appealing baryogenesis mechanism connected to neutrino physics
- various leptogenesis mechanisms:
 - standard leptogenesis: gravitino problem, incompatible with SUSY
 - resonance leptogenesis
 - Dirac leptogenesis
- While there is no model-independent way to test leptogenesis, searches at neutrino experiments (leptonic CPV, neutrino-less double beta decay) can provide supports for/distinguish among the mechanisms
- neutron-antineutron oscillation: complementarity test
 - if observed \Rightarrow low scale leptogenesis scenarios preferred

B violation theory: What did we learn?

- ❖ R. Mohapathra/K. Babu/I. Gogoladze: models exist which give $n\bar{n}$ oscillations within range of improved experiments. Such models tend to possess rather specific structures and also produce signatures at LHC
- ❖ K. Babu: “post-sphaeleron” baryogenesis possibility (which can only be $\Delta B=2$) is NOT ruled out experimentally. Present models tend to make observable LHC predictions.
- ❖ K. Babu/R. Mohapathra: Effective field theory analysis of all $d=9$, $\Delta B=2$ operators in progress (not done before!), might make possible more model-independent statements.
- ❖ M. Chen: “standard” leptogenesis has some problems already! “Resonant” leptogenesis and Dirac leptogenesis also possible (latter since sphaelerons only couple to left-handed components). $N\bar{N}$ possibility is complementary to leptogenesis. Leptogenesis is very difficult to confirm experimentally.

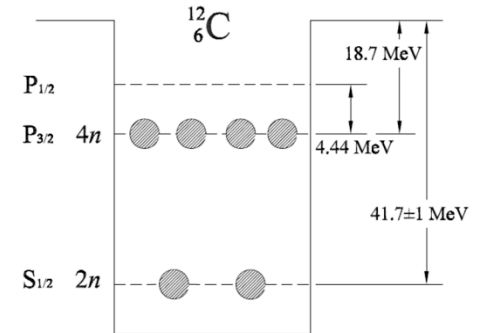
Suppression of $n \rightarrow \bar{n}$ in intranuclear transitions

Neutrons inside nuclei are "free" for the time: $\Delta t \sim \frac{\hbar}{E_{\text{binding}}} \sim \frac{\hbar}{30 \text{ MeV}} \sim 4.5 \times 10^{-22} \text{ s}$

each oscillating with "free" probability $= \left(\frac{\Delta t}{\tau_{n\bar{n}}} \right)^2$

and "experiencing free condition" $N = \frac{1}{\Delta t}$ times per second.

Transition probability per second: $P_A \doteq \frac{1}{\tau_A} = \left(\frac{\Delta t}{\tau_{n\bar{n}}} \right)^2 \times \left(\frac{1}{\Delta t} \right)$



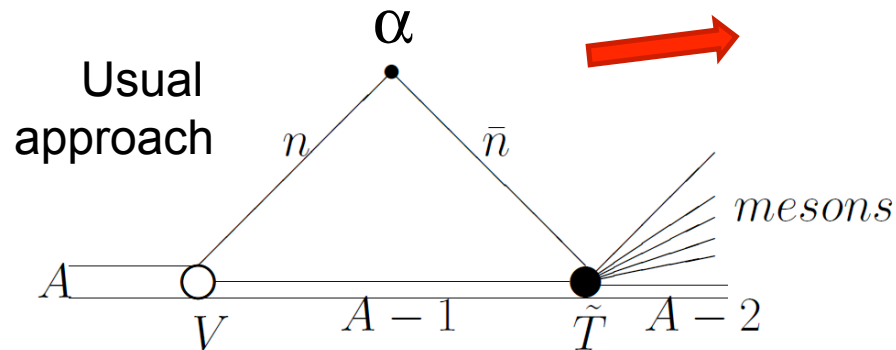
Intranuclear transition (exponential) lifetime:

$$\tau_A = \frac{\tau_{n\bar{n}}^2}{\Delta t} = R \leftrightarrow \tau_{n\bar{n}}^2$$

where $R \sim \frac{1}{\Delta t} \sim 4.5 \leftrightarrow 10^{22} \text{ s}^{-1}$ is "nuclear suppression factor"

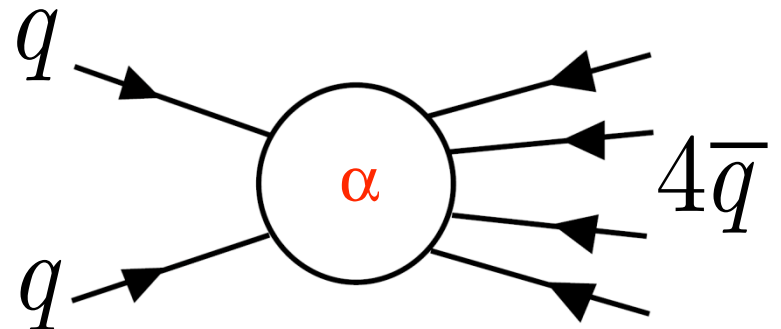
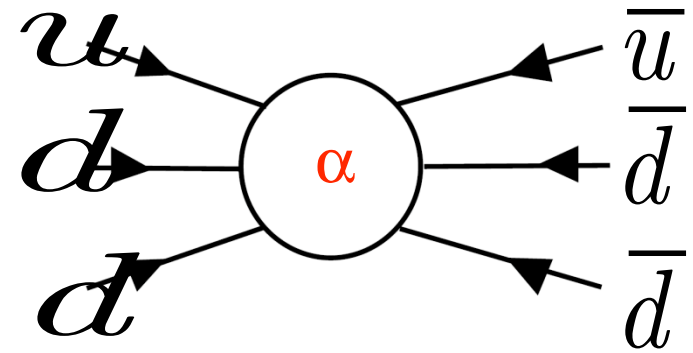
Actual nuclear theory suppression calculations for ^{16}O , ^2D , ^{56}Fe , ^{40}Ar by C. Dover et al; W. Alberico et al; B. Kopeliovich and J. Hufner, and most recently by Friedman and Gal (2008) corrected this rough estimate within a factor of 2

Theoretical nuclear NNbar suppression model is incomplete

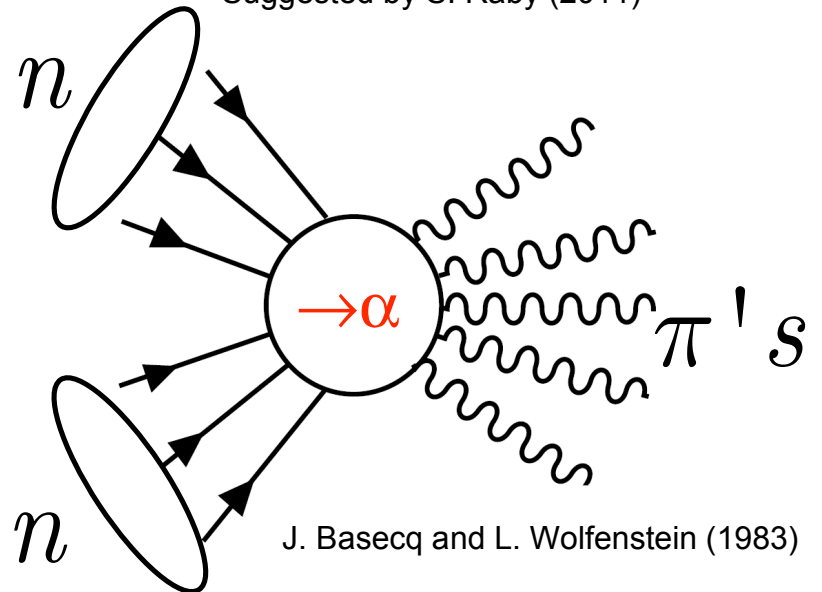


All these processes \rightarrow
include the same amplitude α
and result in the same
indistinguishable final state
(of $\sim 5 \pi s$)

Existing intranuclear NNbar
limits need to be re-evaluated



Suggested by S. Raby (2011)



J. Basecq and L. Wolfenstein (1983)

Estimate

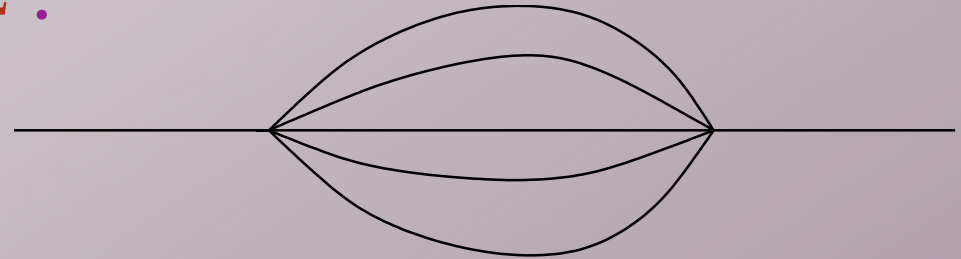
Let us try to use some kind of duality to find a relation between the free $n \leftrightarrow \bar{n}$ oscillation and nuclear stability.

$$\langle \bar{n} | c_{\mathcal{O}}^* \mathcal{O}^\dagger | n \rangle = \epsilon \bar{u}_{\bar{n}}^c \gamma_5 u_n \quad |\epsilon| = \frac{\hbar}{\tau_{n\bar{n}}}$$

where \mathcal{O}^\dagger decreases B , $\Delta B = 2$.

Operator product expansion

$$\int d^4x e^{iqx} T\{\mathcal{O}(x)\mathcal{O}^\dagger(0)\} = c_q \bar{q}q + \dots$$

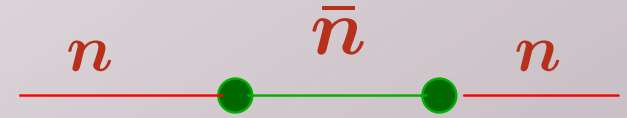


The average over a nucleus A gives its lifetime τ_A

$$2|c_{\mathcal{O}}|^2 \text{Im} \int d^4x \langle A | T\{\mathcal{O}(x)\mathcal{O}^\dagger(0)\} | A \rangle = \frac{\hbar}{\tau_A}$$

The average over neutron state

$$|c_{\mathcal{O}}|^2 \int d^4x e^{iqx} \langle n | T \{ \mathcal{O}(x) \mathcal{O}^\dagger(0) \} | n \rangle \sim \frac{|\epsilon|^2}{\Delta}$$



where Euclidian $q \sim \Delta$ is a relevant hadronic duality scale.

Taking $\langle A | \bar{q}q | A \rangle \sim A \langle n | \bar{q}q | n \rangle$ for the leading OPE term we get

$$\tau_A = R \tau_{n\bar{n},n}^2 \quad R = \frac{\Delta}{A\hbar}$$

For ^{16}O and an educated guess for $\Delta = 0.5 \text{ GeV}$

$$R = 4.7 \times 10^{22} \text{ s}^{-1}$$

what is close to the result obtained by Friedman, Gal (2008).

The inclusive approach does include all the mechanisms.

PROTON DECAY

- ✱ Proton is a topological non-trivial configuration of the pion field (Skyrmion)
- ✱ Decay of the proton is protected by topology
- ✱ Hybrid Skyrmion/bag model decay possible but exponentially suppressed due to tunneling (instanton)

DISCUSSION

- ✱ We calculated hadronic matrix elements including non-perturbative QCD effects resulting in suppression.
- ✱ This suppression can be sizeable.
- ✱ Drawback not a very stable calculation due to bag size.

Where Lattice Can Help

- ♦ Is BSM running non-perturbative?
 - Model-dependent (assume pert. models for now)
- ♦ Is QCD running non-perturbative?
 - Should be checked (pert. running reasonable)
- ♦ What is neutron-antineutron matrix element?
 - Inherently non-perturbative question
- ♦ What is effect in nuclei?
 - Very interesting, VERY hard question

Future Outlook

Currently in progress:

- ✦ Independent analysis checks
- ✦ $L = 20, 390 \text{ MeV}$ pions
- ✦ $L = 32, 240 \text{ MeV}$ pions

Feasible in the next year or two:

- ✦ Physical Point Calculation
- ✦ Chiral Fermion Calculation



NNbar suppression factor in nuclei: theory developments

A. Vainshtein: operator product expansion calculation in progress (with B. Kopeliovich) will implicitly include all processes and give independent estimate of size and error of Gal calculation.

M. Buchoff: lattice calculation of nnbar transition matrix element in progress, special structure of nnbar operator makes it possible, should make possible quantitative connection between nnbar limit and energy scale

M. Stavenga: Skyrme calculation of extra suppression of B violation from chiral dynamics?

ALSO (Vainshtein): $\Delta B=2$ in nuclei can also come from “di-proton decay”, How does this affect limits from nnbar in nuclei?

Vacuum N-Nbar transformation from bound neutrons:

Best result so far from Super-K in Oxygen-16

$$\tau_{^{16}\text{O}} > 1.89 \leftrightarrow 10^{32} \text{ yr} \quad (90\% \text{ CL})$$

\Re 24 observed candidates;
24.1 exp. background

$$\tau_{\text{nucl}} = R \times \tau_{n\bar{n} \text{ free}}^2$$

if $R_{^{16}\text{O}} = 5 \cdot 10^{22} \text{ s}^{-1}$ (from Friedman and Gal 2008)

$$\Rightarrow \tau(\text{from bound}) > 3.5 \times 10^8 \text{ s} \quad \text{or} \quad \alpha < 2 \times 10^{-24} \text{ eV}$$

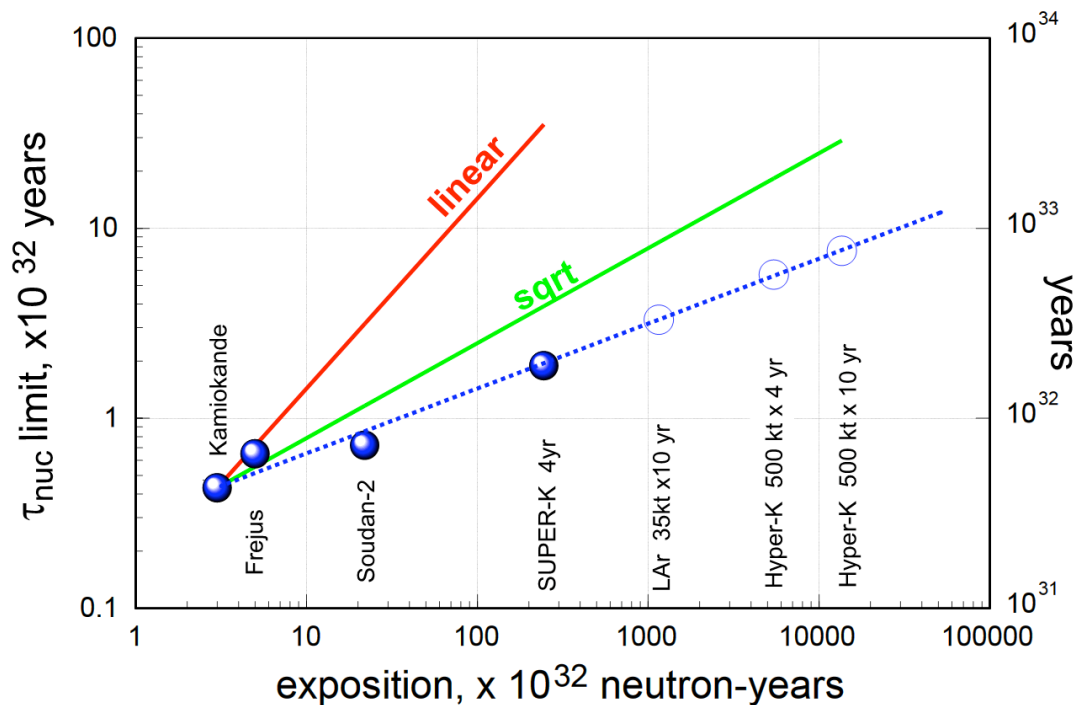
\leftrightarrow 16 times higher than
sensitivity of ILL expt.

ILL limit (1994) for free neutrons: $\tau_{n\bar{n}} > 0.86 \times 10^8 \text{ s}$

Bound neutron N-Nbar search experiments

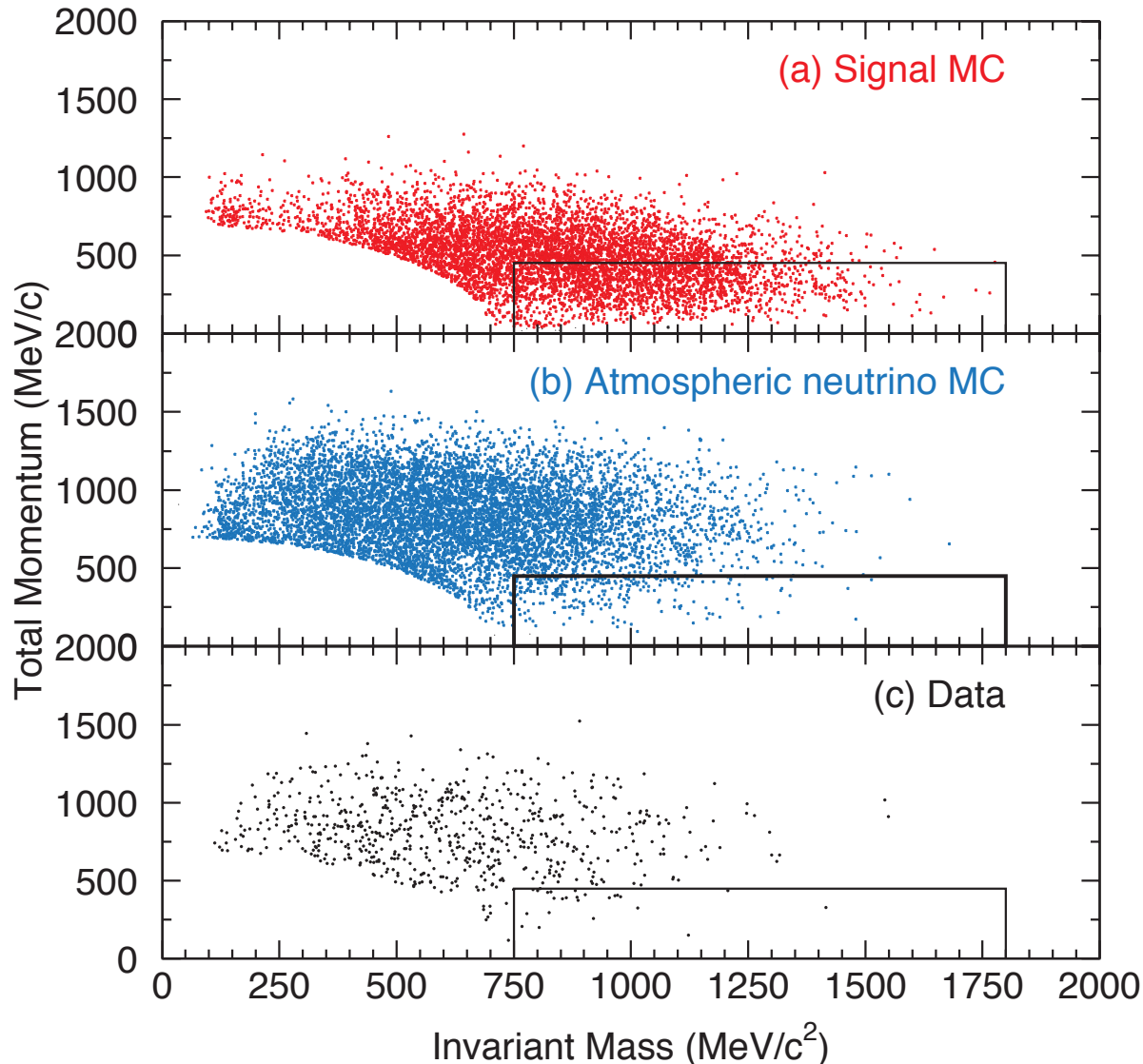
Experiment	Year	A	n-year (10^{32})	Det. eff.	Candid.	Bkgr.	τ_{nucl} , yr (90%)
Kamiokande	1986	O	3.0	33%	0	0.9/yr	$>0.43 \times 10^{32}$
Frejus	1990	Fe	5.0	30%	0	4	$>0.65 \times 10^{32}$
Soudan-2	2002	Fe	21.9	18%	5	4.5	$>0.72 \times 10^{32}$
SNO *	2010	D	0.54	41%	2	4.75	$>0.301 \times 10^{32}$
Super-K	2011	O	245	12.1%	24	24.1	$>1.89 \times 10^{32}$

* Preliminary



- From Kamiokande to Super-K atmospheric ν background is about the same ~ 2.5 /kt/yr.
- Large D₂O, Fe, H₂O detectors are dominated by backgrounds; LAr detectors are unexplored
- Observed improvement is weaker than SQRT due to irreducible background and uncertainties of efficiency and background.
- Still possible to improve a limit but impossible to claim a discovery.

Super-Kamiokande Result



12 % detection efficiency
 sys. uncertainty 23%
 (mostly intranuclear scattering)

24.1 background events
 ν osc. effects are included
 sys. uncertainty 24%
 (mostly flux, cross sections)

24 candidates

$$T_{bound} > 1.89 \times 10^{32} \text{ years}$$

$$\tau_{free} = \sqrt{\frac{T_{bound}}{1 \times 10^{23} \text{ s}^{-1}}} \\ = 2.4 \times 10^8 \text{ s}$$

Liquid Argon TPC

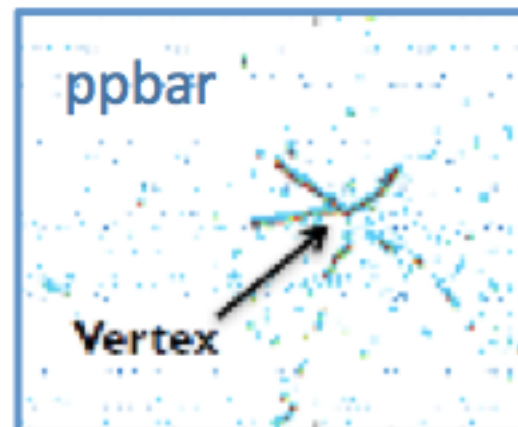
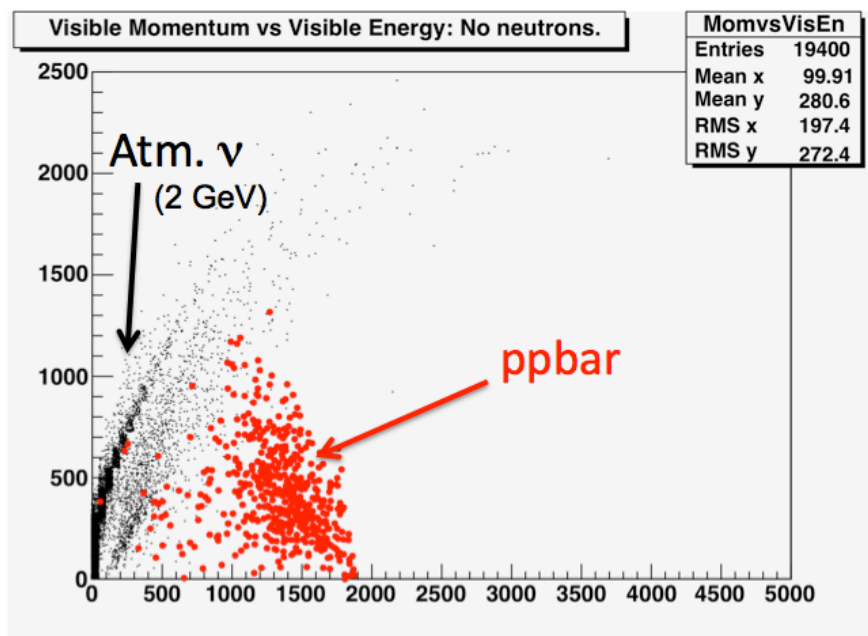
Compared to Iron Calorimeters:

- can do better than requiring $n_{\text{ch}} \geq 4$

Compared to WC

- can resolve recoil proton, charged current lepton

Potentially big gains
in efficiency and
BG rejection!



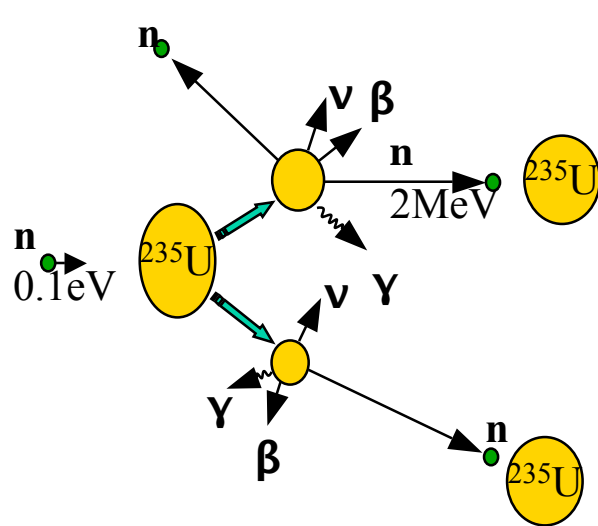
Good discrimination at least at truth level.

G. Karagiorgi, LBNE-docdb-5645

Observations

- ❖ Proton decay detectors have a long history of studying \bar{n} . Usual qualities apply:
large mass, high efficiency, low background
- ❖ Analyses have been fairly crude so far.
No modern MVA techniques. High background rate in water cherenkov is daunting.
- ❖ LAr TPC, even one as small as LBNE/10 kton should do very well. Let's study!

“Slow” Neutrons: MeV to neV

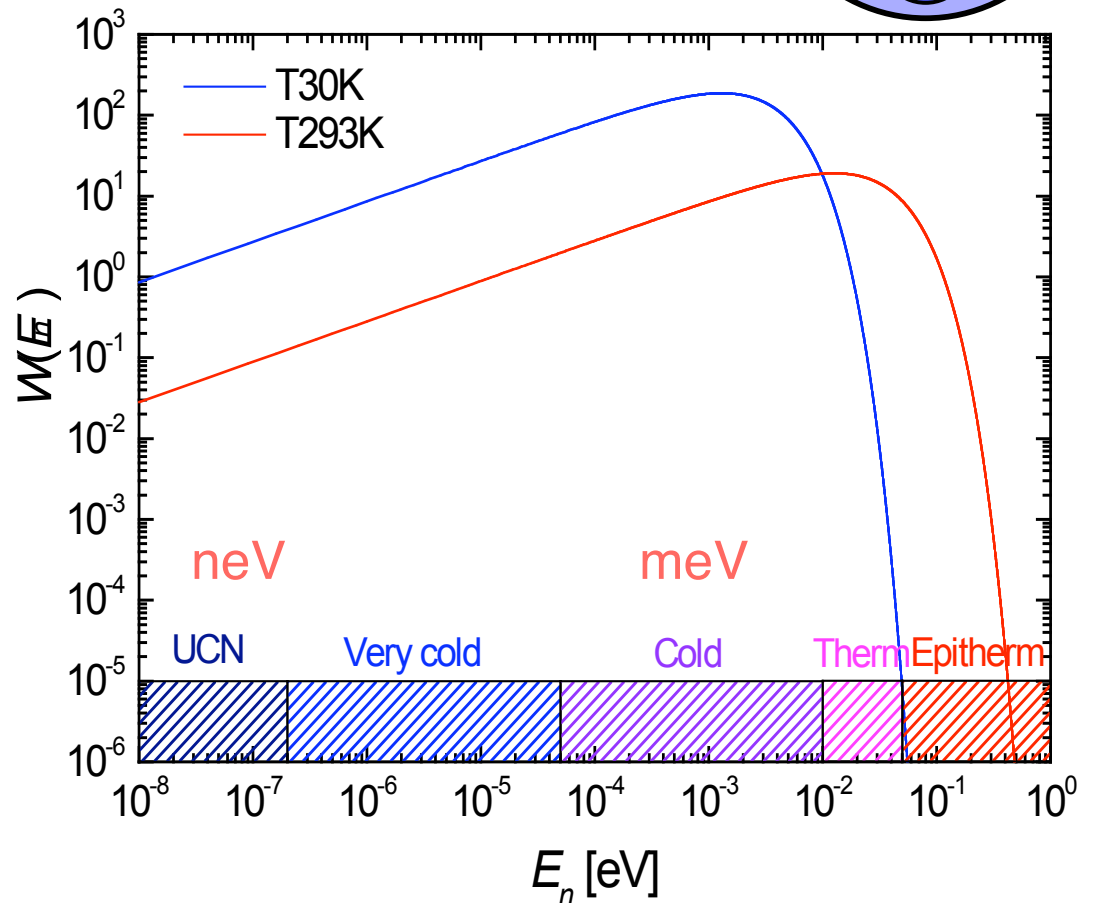
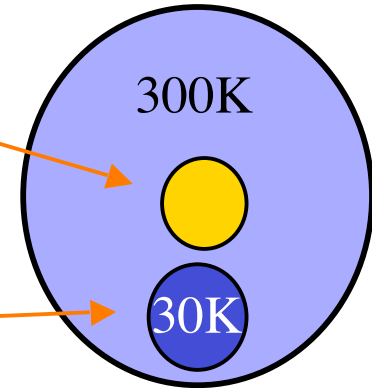


~MeV neutrons from fission or spallation, thermalized in ~ 20 collisions in ~ 100 μ s

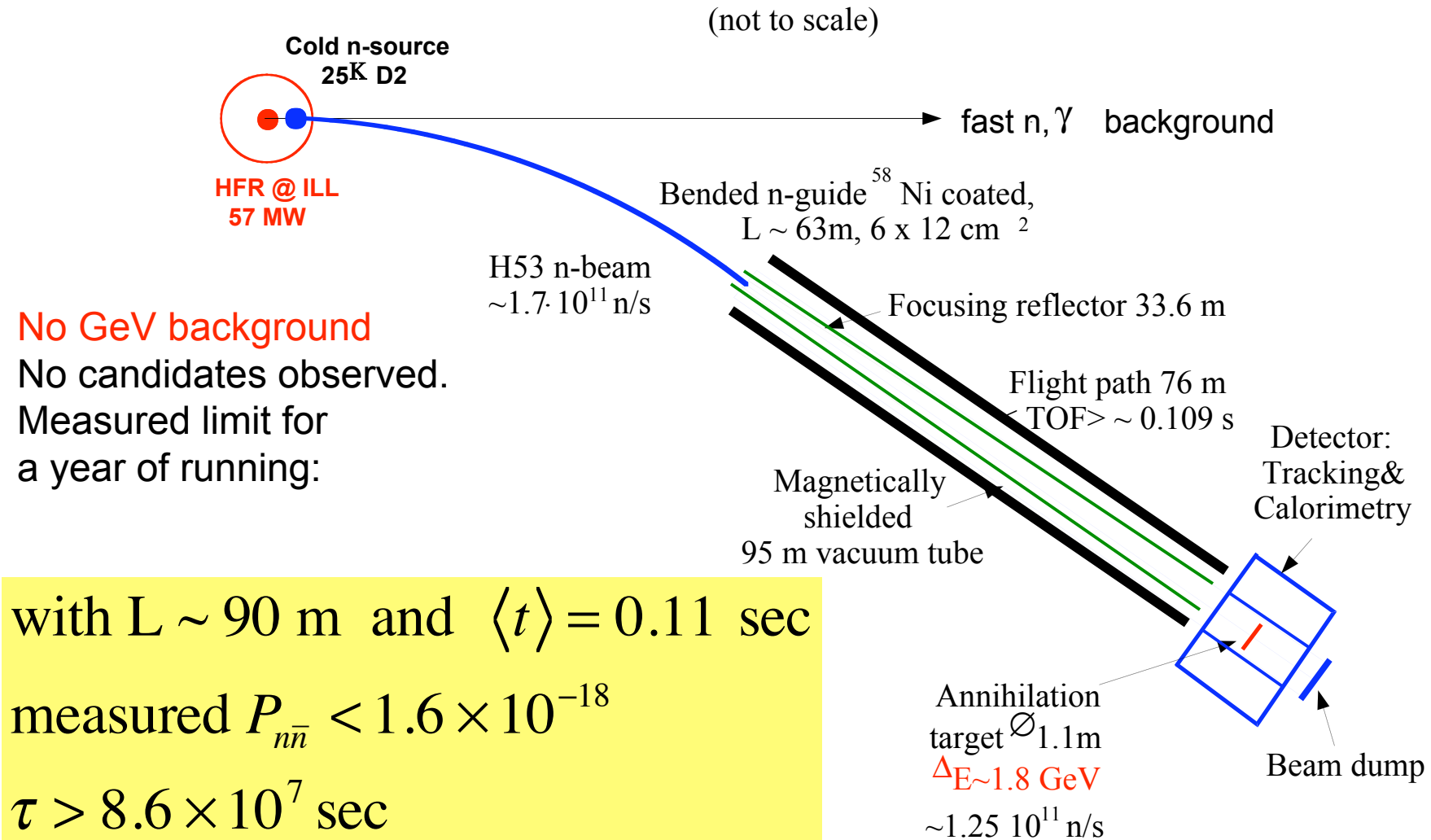
T (K)	E (meV)	λ (Å)	v (m/sec)
300	25	1.6	2200
20	2	6.4	550

Nuclear reactor/
Spallation source

Neutron Moderator
(LH2, LD2)



N-Nbar search at ILL (Heidelberg-ILL-Padova-Pavia)



Baldo-Ceolin M. et al., Z. Phys. C63,409 (1994).

Quasifree Condition: B Shielding and Vacuum

$\mu B t \ll \hbar$ ILL achieved $|B| < 10$ nT over 1m diameter, 80 m beam, one layer 1mm shield in SS vacuum tank, 1% reduction in oscillation efficiency (Bitter et al, NIM A309, 521 (1991). For new experiment need $|B| < \sim 1$ nT

If nnbar candidate signal seen, easy to “turn it off” by increasing B

$V_{\text{opt}} t \ll \hbar$:

Need vacuum to eliminate neutron-antineutron optical potential difference.

$P < 10^{-5}$ Pa is good enough, much less stringent than LIGO

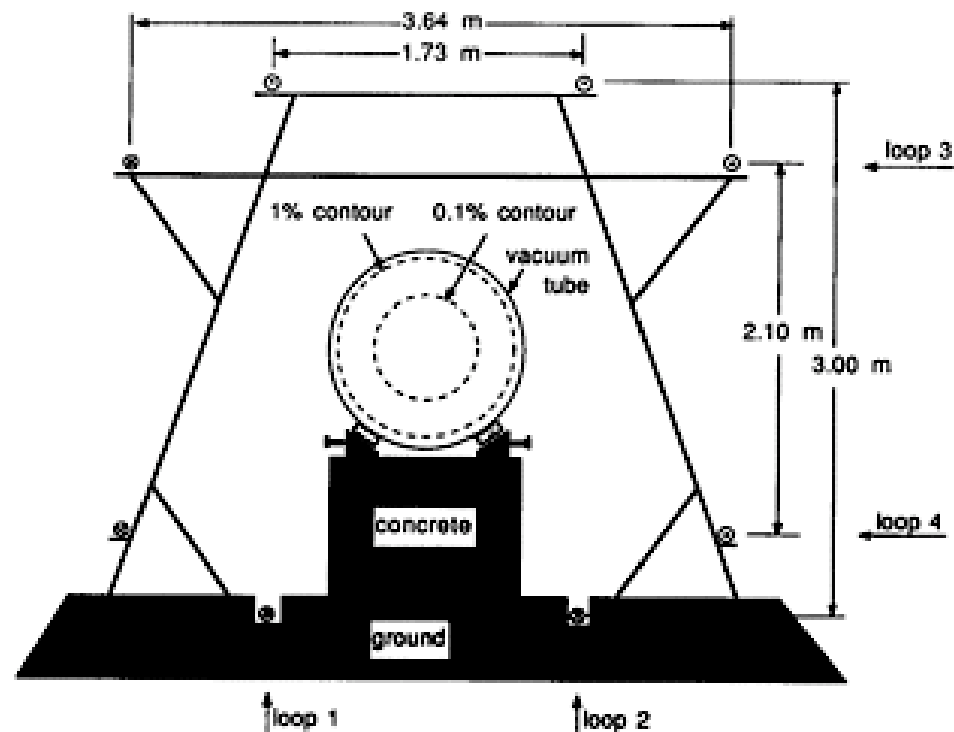


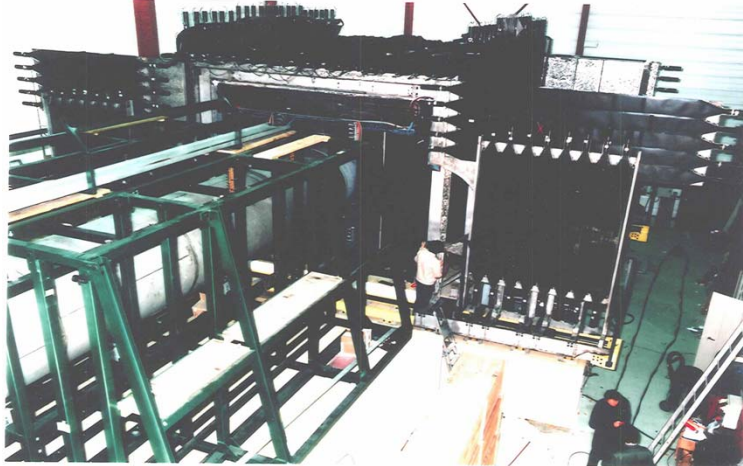
Fig. 10. The transverse field compensation system. Loops 1 and 2 are under 49 A current and compensate the horizontal field component; loops 3 and 4 are under 120 A current and compensate the vertical field component.

2. ILL n-nbar beam line

Cold
neutrons →

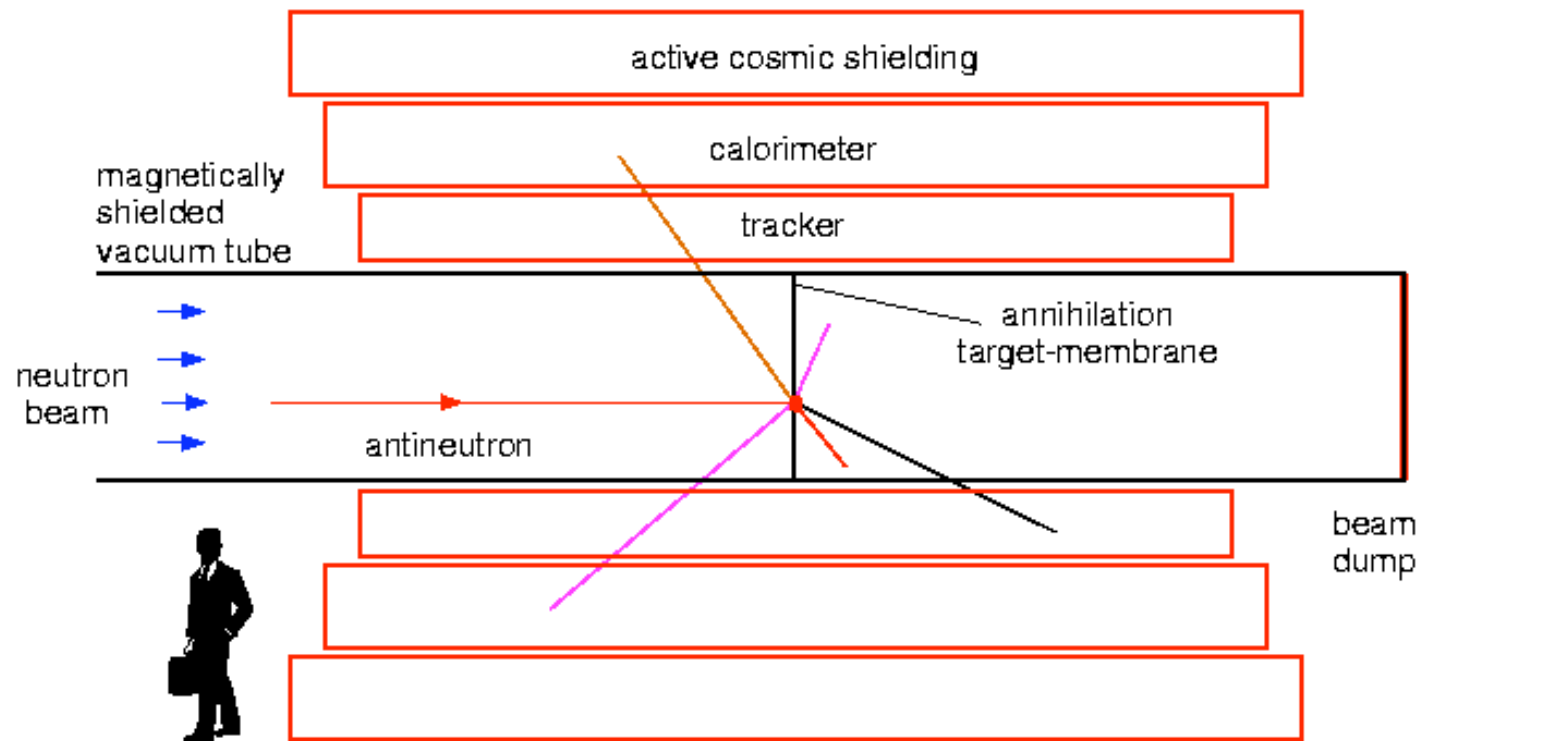


Annihilation detector



Beam
stop

The conceptual scheme of antineutron detector



$$\bar{n} + A \rightarrow \langle 5 \rangle \text{ pions} \quad (1.8 \text{ GeV})$$

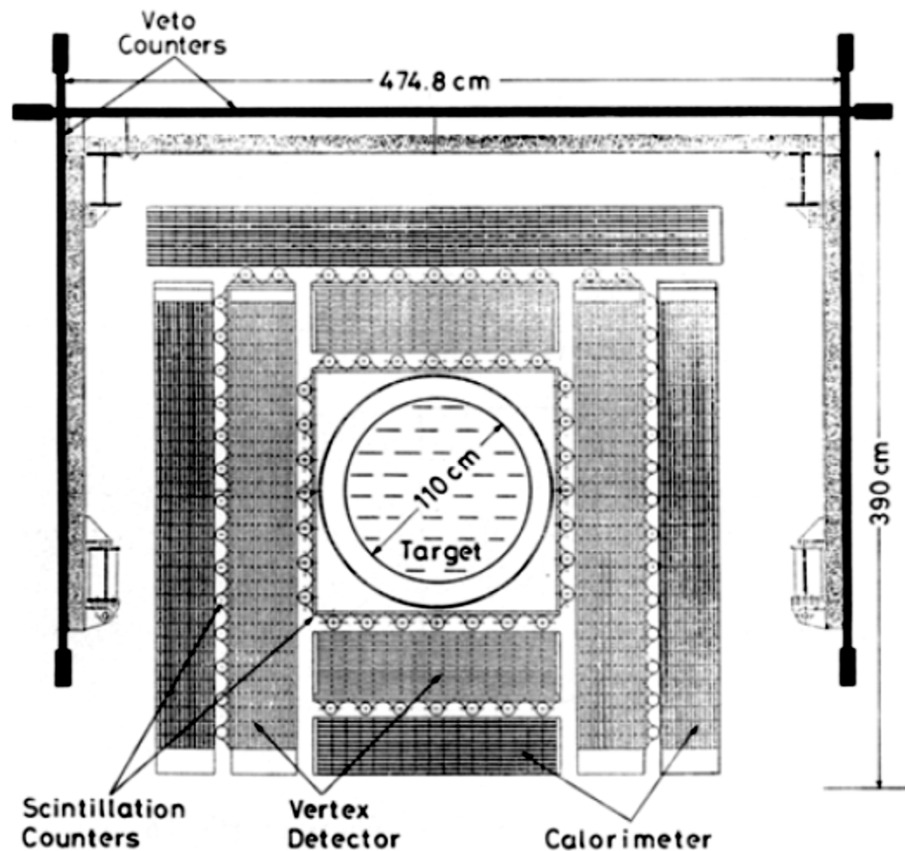
Annihilation target: $\sim 100\mu$ thick Carbon film

$\sigma_{\text{annihilation}} \sim 4 \text{ Kb}$ $\sigma_{\text{nC capture}} \sim 4 \text{ mb}$

vertex precisely defined. No background was observed

Annihilation detector (INFN Padova and Pavia)

1. Inner Vertex Detector: 10 layers of Limited Streamer Tubes (LST), 0.3 g/cm^3 , Vertex $\pm 4 \text{ cm}$
2. Outer Calorimeter: 12 layers of LST interleaved with Pb/Al planes
3. Timing: Inner and outer planes of Plastic Scintillators (PSc), 700 ps,
4. Cosmic ray rejection with 95 m^2 outmost layer of PSc, separated by 10 cm Pb.



60 000 electronic channels
Overall $n\bar{n}$ detection
efficiency $52 \pm 2\%$.

Explosion-proof gas mixture

Fig. 5. The $n\bar{n}$ annihilation detector (cross sectional view)

$n\bar{n}$ at ILL

Information from D. Dubbers, based on ILL Experiment

The <10 nT stated limit was conservative. ~ 1 nT should be achievable with a very similar shielding approach. Need to also worry about 60 Hz

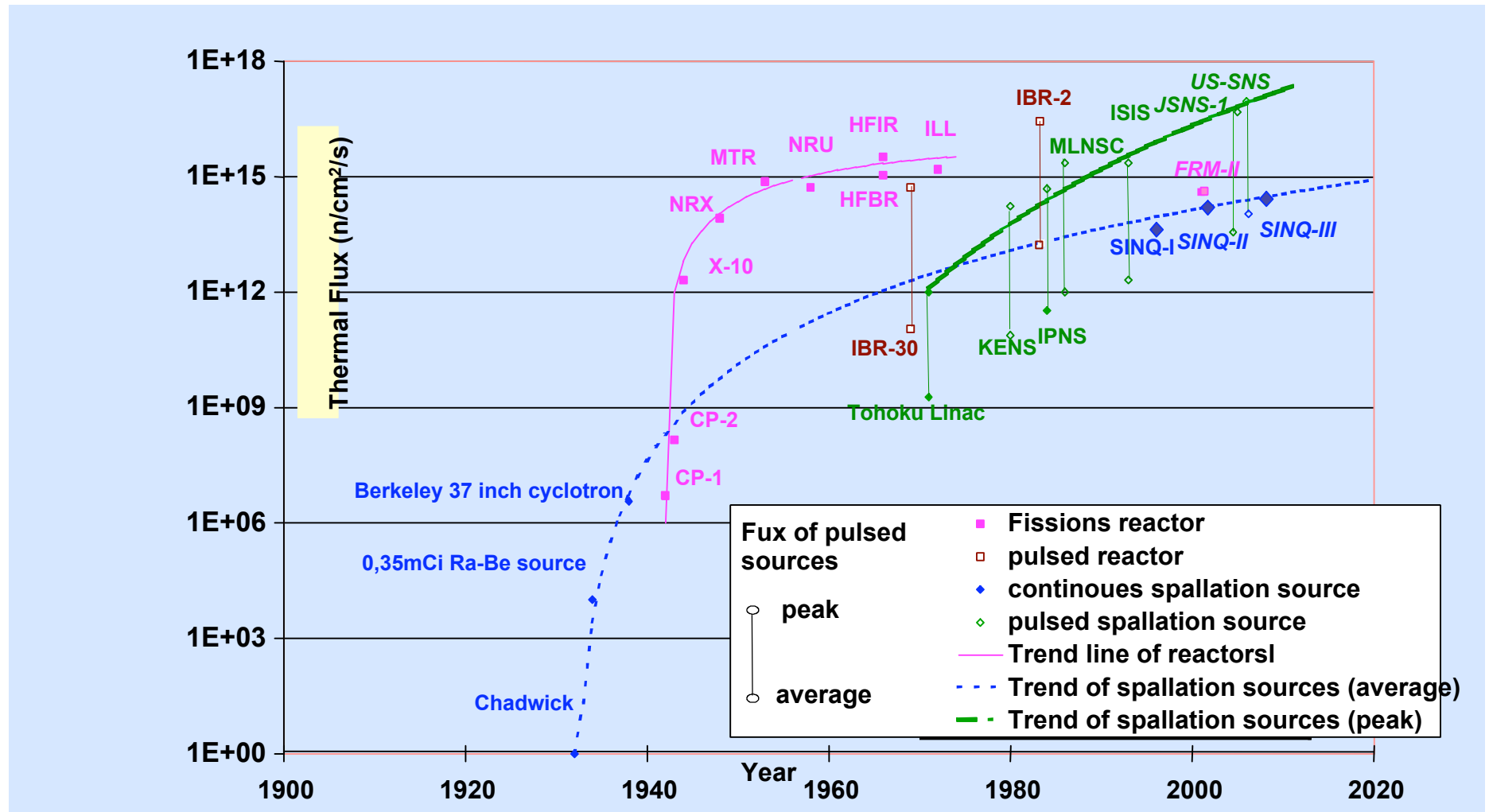
Vertex resolution of ILL nnbar detector was very coarse (~ 5 cm) compared to annihilation target thickness (~ 100 microns). Lots of room for even further background reduction.

Neutron backgrounds from slow neutron absorption/scattering on annihilation target can be (and needs to be) improved in new experiment to reduce tracker deadtime from MeV capture gammas

Vacuum chamber/B shielding of experiment still exists at ILL

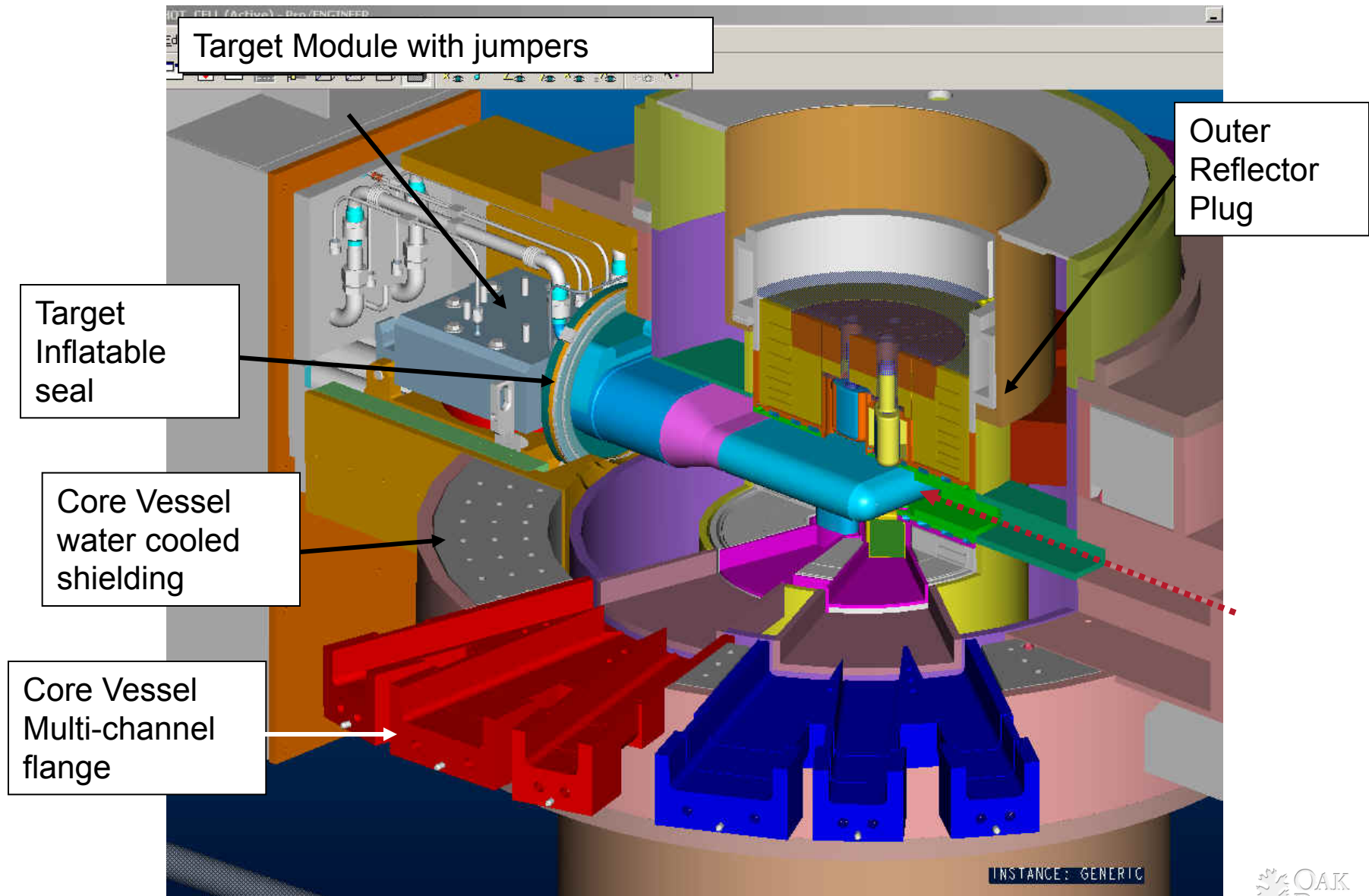
How to Improve the Experiment? Not so Easy.

Max neutron flux/brightness: ~unchanged for ~4 decades



Neutron flux is increasing only slowly with time R. Eichler, PSI

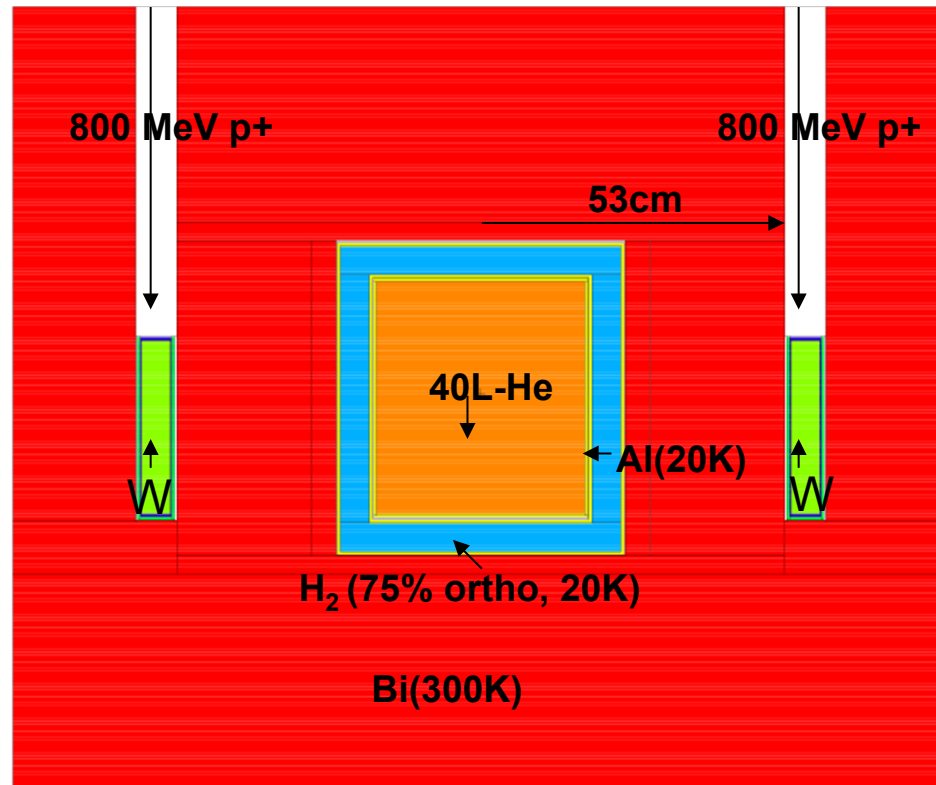
Target Region Within Core Vessel



Summary

- **The SNS is operating at a very high level of reliability and at times power levels $> 1\text{MW}$.**
- **Development of high powered targets based on the SNS experience can be accomplished.**
- **Cost savings are possible based on the SNS data.**
- **Experienced personnel are available to help develop these high powered targets.**

Inverse cylindrical geometry (1)



$6.6 \cdot 10^7$ UCN/s/100mA

Heat load @ 100mA \equiv 80KW

Total heat: 27.4 W

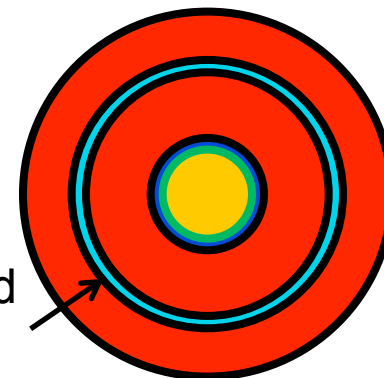
Neutron heat: 17.2 W

Photon heat: 9.6 W

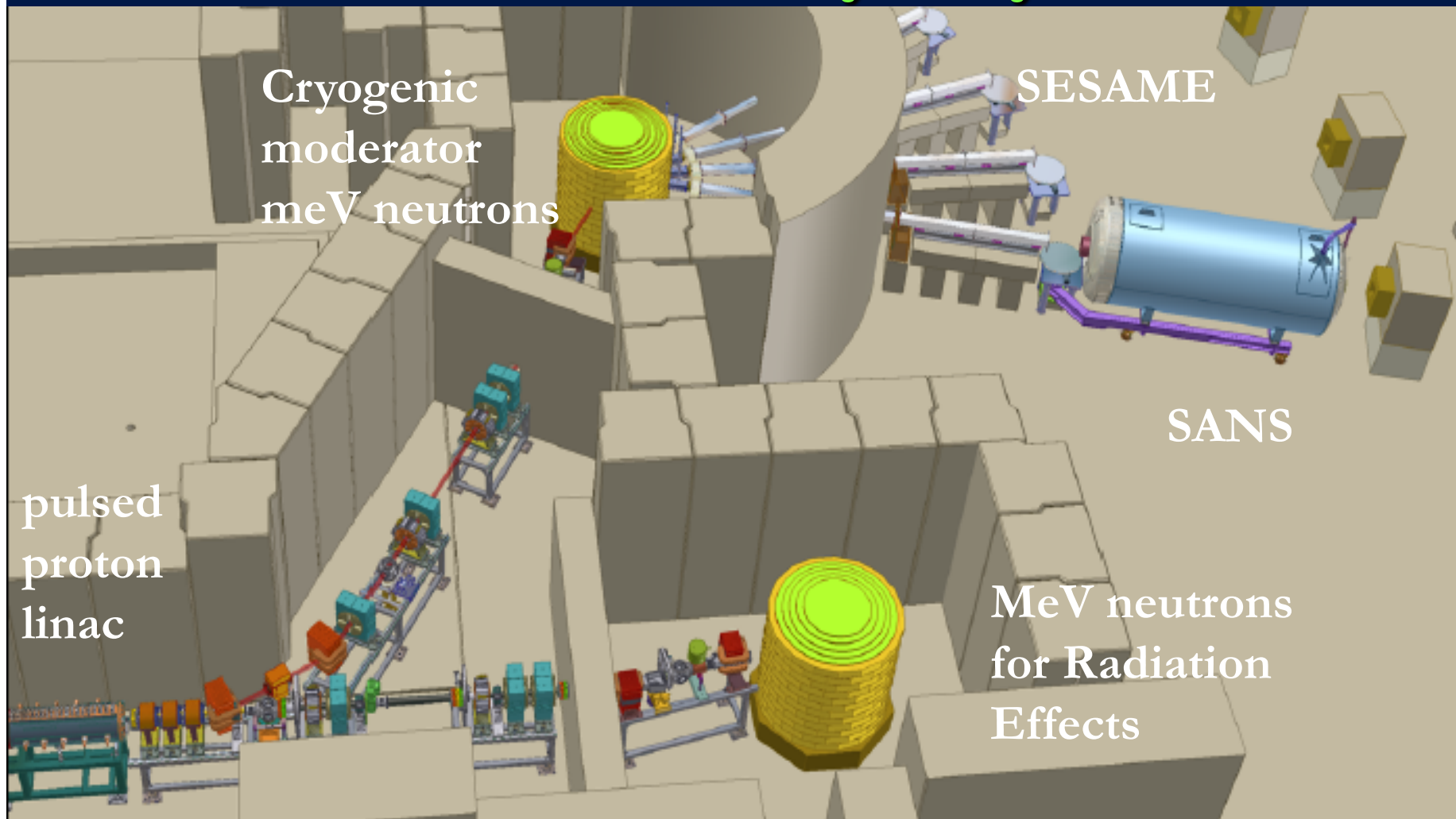
Proton heat: 0.6 W

$2.4 \cdot 10^8$ UCN/s/100W (heat in the He)

Cylindrical proton target (beam rastered around circumference)



LENS Facility Layout



Designed/built/characterized by graduate students
Local user program in operation

~1MW Slow Neutron Source @Project X?

G. Greene: rough scaling from SNS+ straight guide->~1/4 ILL possible

T. Gabriel: project X source would be less \$\$\$ than SNS, many benefits from SNS experience and ongoing ESS design

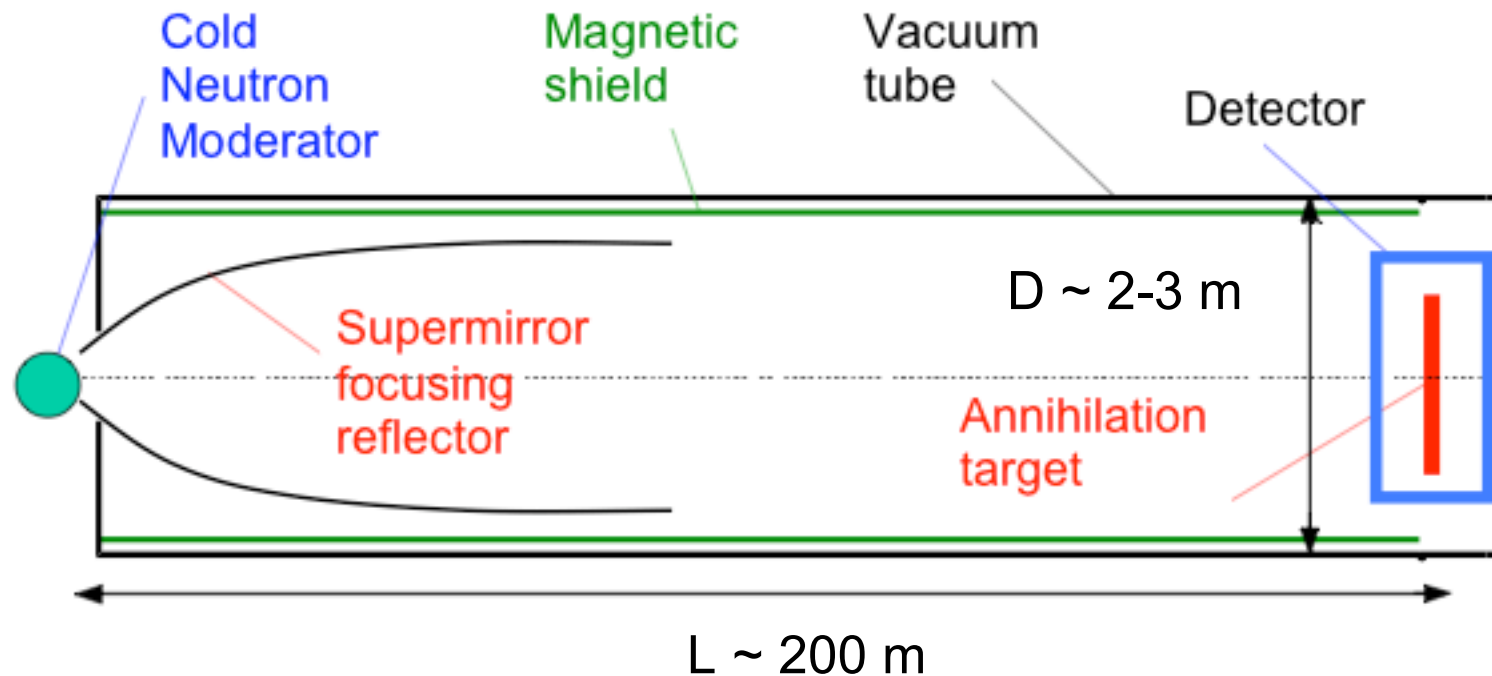
G. Muhrer: MCNP/vetted design for cold source with high kappa superfluid helium exists.

C. Liu (for D. Baxter): LENS neutron source at IU can be used to evaluate cold n moderator improvements (grooved moderators, nanoparticle reflectors,...)

Better Free Neutron Experiment (Horizontal beam shown: vertical possible)

need slow neutrons from high flux source, access of neutron focusing reflector to cold source, free flight path of $\sim 200\text{m}$

Improvement on ILL experiment by factor of ~ 1000 in transition probability is possible with existing n optics technology (see G. Greene talk)



Possible improvements in sensitivity (Nt^2)

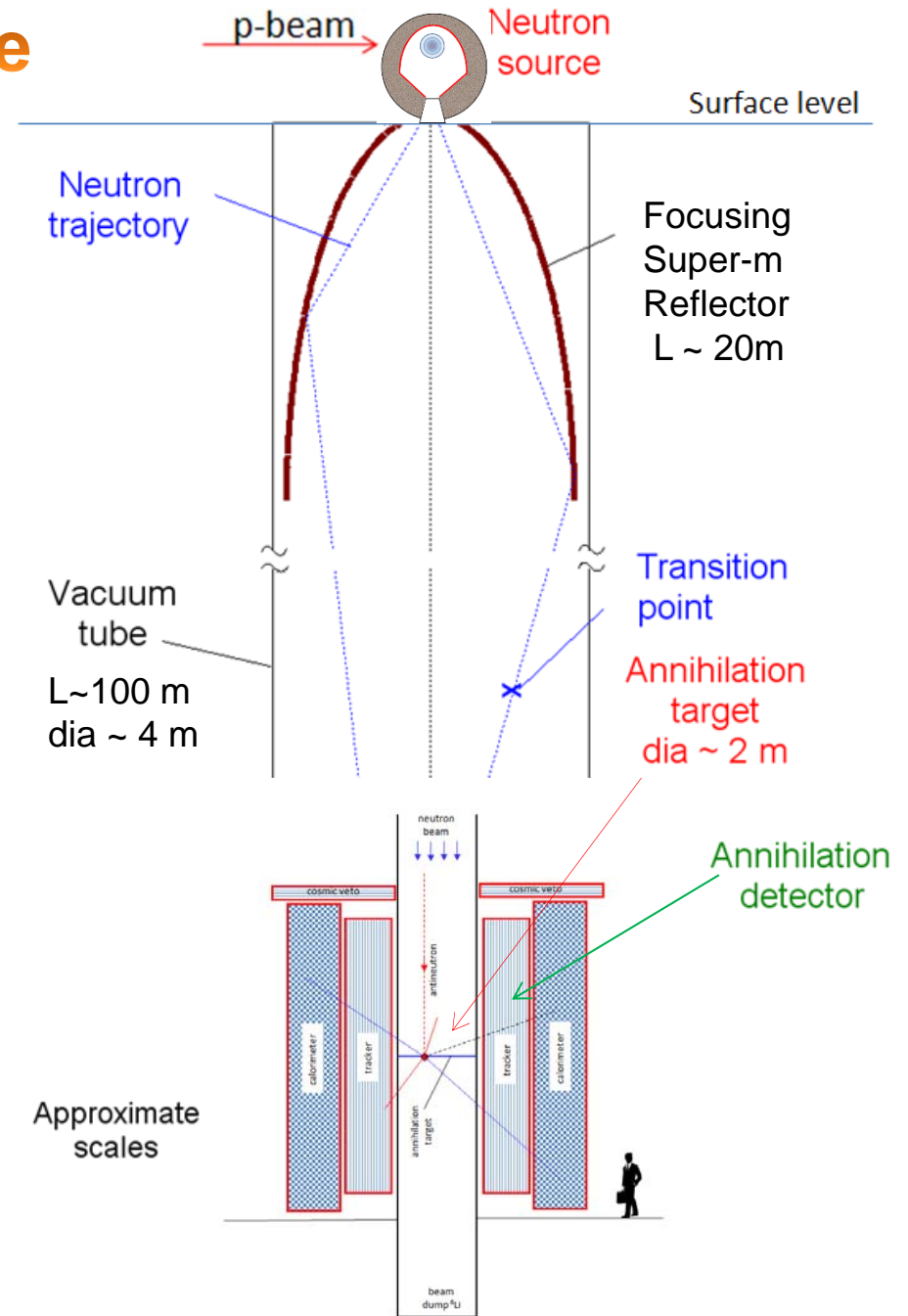
- Intrinsic source brightness (assume 1MW) x 1/4
- Colder moderator (gain goes as λ^2) x 2
- Coupling to experiment x 2
- Larger moderator face ($30 \times 30 \text{ cm}^2$ vs $6 \times 12 \text{ cm}^2$) x 12
- Use “high-m” neutron reflector (assume $m=6$) x 36
- Longer experiment (200m vs 76m gain $\sim L^2$) x 7

Estimated Sensitivity Gain $\sim 3 \times 10^3$

Take away message: A substantial improvement is possible with only straightforward extension of existing technology

Top-Down vertical scheme

- Can combine most of improvements;
- CW or pulsed;
- Max UCN (<10 m/s) enrichment will be most advantageous;
- Cold and VCN are also used;
- Ultimate combination of all improvements should boost the sensitivity by factor $> 1,000$ \times times several years of operation

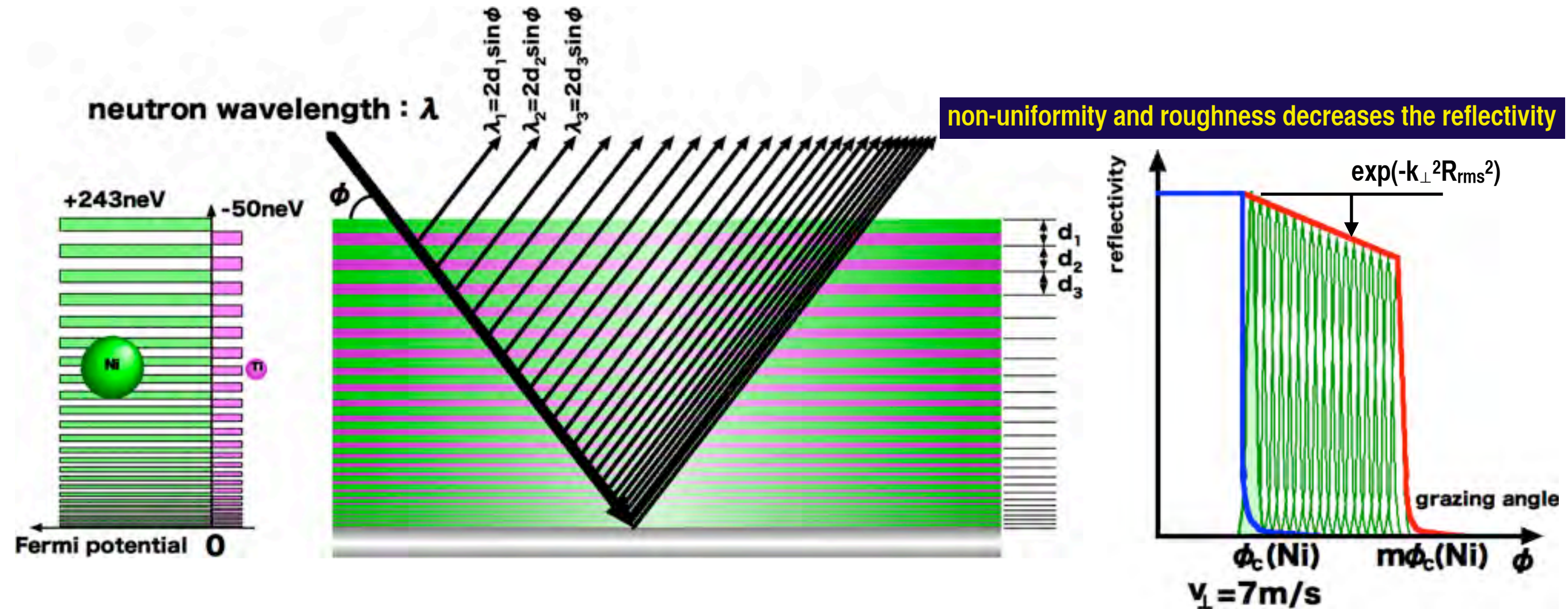


$$h = v_0 t + \frac{1}{2} g t^2$$

$$105 \text{ m} = 100 \text{ m/s} \cdot 1 \text{ s} + 4.9 \text{ m/s}^2 \cdot 1^2 \text{ s}^2$$

$$105 \text{ m} = 10 \text{ m/s} \cdot 3.7 \text{ s} + 4.9 \text{ m/s}^2 \cdot 3.7^2 \text{ s}^2$$

Supermirror



$$m = \phi_c / \phi_c(\text{Ni}) = v_c(\text{Ni}) / v_c$$

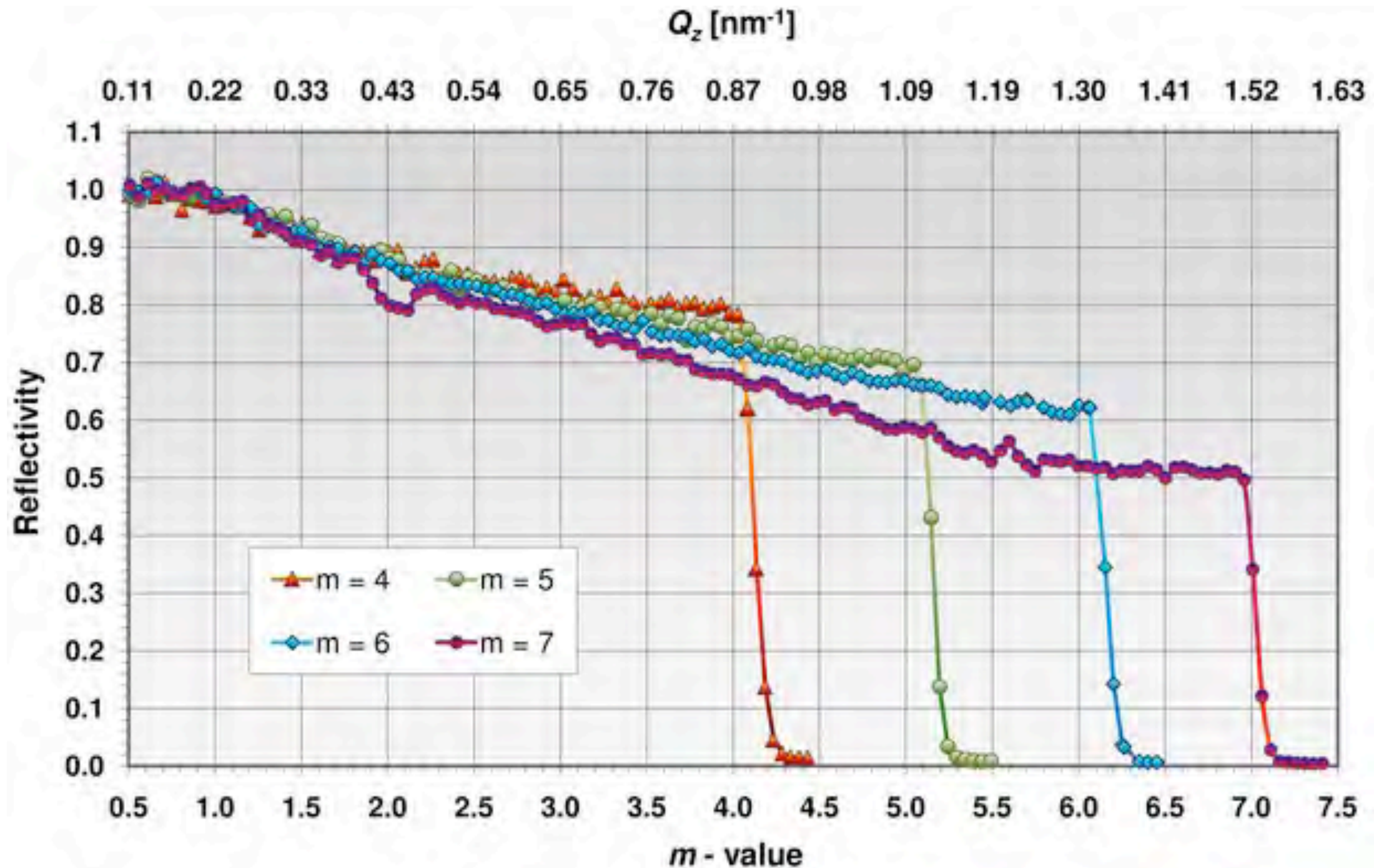
$$\phi_c(\text{Ni}) / \lambda = 1.7 \text{ mrad/\AA}$$

$$v_{\perp}(\text{Ni}) = 7 \text{ m/s}$$

m=4-7 Supermirrors

<http://www.swissneutronics.ch/>

Supermirror: commercially available up to $m=7$ ($v_{\perp}=50\text{m/s}$)



Summary

Multilayer mirrors enhances the figure-of-merit of n - \bar{n} experiments.

Multilayer fabrication technology was remarkably improved in the past decade.

monochromatic reflectors $m \leq 10$

Focusing of cold neutrons in vertical flight path

supermirrors $m \leq 7$

Confinement of VCN

Enhancement of VCN intensity

substrateless supermirrors $m \leq 5$

Enhancement of VCN intensity



Supermirror Optics

G. Greene: greatest single contributor to possibility of improved free neutron experiment

H. Shimizu: $m=10$ multilayer monochromators exist, $m=7$ supermirrors, exist, radiation damage can be handled using SM coating on metal, research on H and D-doped diamond-like carbon mirrors in progress

H. Shimizu: Nagoya U active x-ray mirror manufacturing group exists, available ~2015 for new project

Groups in India



- ❑ During May, 2011, a short workshop was organized by Dr. Amlan Ray in VECC, Kolkata on N-Nbar oscillation studies
- ❑ Several experts from USA participated in this event
- ❑ A group from VECC (Kolkata) led by Dr. Ray had a few discussions with the Nuclear and Particle physics groups at SINP (Kolkata)
- ❑ The 2 institutes jointly show interest in joining an activity on N-Nbar oscillation studies
 - P. Das, A. Ray, A.K. Sikdar at VECC
 - S. Banerjee, S. Bhattacharya, S. Chattopadhyay at SINP

Free neutron $n\bar{n}$ search: relation with other project X ideas?

Technical:

B. Filippone: both $n\bar{n}$ and (one version of) $n\text{EDM}$ can use bright slow neutron source: might one source feed both?

(someone in tracker session): detectors for $\mu 2e$ experiment and kaon experiments share neutron-induced background issues with $n\bar{n}$ detector

Scientific:

$N\bar{n}$ improvements squeeze post-sphaeleron baryogenesis. EDM experiments squeeze sphaeleron+EW-scale BSM physics. Do null measurements in both areas at Project X/elsewhere leave leptogenesis by default as the last viable baryogenesis mechanism?

3 Questions

1. *How much better well could we do at Project X?*

MUCH BETTER... BUT NEED DETAILED SIMULATIONS

2. *What would it cost?*

NEED PRELIMINARY ENGINEERING

3. *Is it worth doing?*

NEED ANSWERS TO 1. & 2. PLUS THEORY

NNbar and Project X: What do we need (what will we have?) by Snowmass?

Theory:

sharper understanding of nnbar in nuclei

EFT analysis of all $\Delta B=2$ operators involving standard model fields

(preliminary) lattice calculations of nnbar matrix element

Experiment (underground detectors):

Calculation of $\Delta B=2$ reach for underground liquid Ar detectors

Experiment (free neutrons):

Sensitivity/\$\$\$ ratios for likely options

NNbar Summary

New physics beyond the SM can be discovered by NNbar search

Improvement in free neutron oscillation probability of a factor of $\sim 1,000$ is possible

If discovered:

- $n \rightarrow \bar{n}$ observation would violate B-L by 2 units, establish a new force of nature, illuminate beyond SM physics, and may help to understand matter-antimatter asymmetry of universe

If NOT discovered:

- will set a new limit on the stability of “normal” matter via antimatter transformation channel. Will constrain some scenarios for B-L violation and “post-sphaeleron” baryogenesis

Summary

New physics beyond the Standard Model can be discovered by $\bar{N}N$ search

Experiments with free neutrons possess very low backgrounds (sharp vertex localization): ILL experiment observed no background. Interpretation of result is independent of nuclear models. Any positive observation can be turned off experimentally with the application of a small magnetic field.

Sensitivity of free neutron experiment for $\bar{N}N$ transition rate can be improved by factor of ~ 1000 using existing technology [Combination of improvements in neutron optics technology, longer observation time, and larger-scale experiment]. Further improvements in a free neutron experiment can come from neutron optics technology development.

US high-energy intensity frontier complex could in principle provide the type of dedicated source of slow neutrons needed for $\bar{N}N$ experiment.